Low-Temperature Raman Study on ThSiO₄ Single Crystal and Related Infrared Spectra at Room Temperature

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Before analyzing the optical spectra of some M^{4+} ions diluted in ThSiO₄, Raman and infrared studies were undertaken on a single crystal of thorite. Indeed a strong coupling between the electronic transitions within the f band and the vibrations of the host lattice is expected. All the active modes of vibration at k = 0 of ThSiO₄ were assigned in the D_{4h} point group. Raman scattering spectra recorded between 300 and 4.2 K have shown that the crystalline structure of thorite does not change down to 4.2 K. @ 1986 Academic Press. Inc.

Introduction

Thorium silicate, ThSiO₄, is an attractive crystalline host for the study of the electronic structure of M^{4+} actinide ions by optical spectroscopy and these doped material may have some possibility in nonlinear optics and quantum electronics. The tetragonal transparent crystal stabilizes the 4+ oxidation state of the actinide ions substituted into the D_{2d} symmetry site of Th⁴⁺. Due to the highest radius of Th⁴⁺ among the M^{4+} ions of the actinide series, the substitution can occur without distorting the neighboring environment.

For a complete analysis of the electronic and vibronic spectra, it is necessary to know the energy of the vibration modes of the ThSiO₄ lattice. Moreover, the optical spectra being recorded at 4.2 K we have to be sure that the crystalline structure of ThSiO₄ remains the same at low temperature. Therefore, Raman studies were performed at different temperatures varying from 300 to 4.2 K, as well as infrared absorption at room temperature.

Until recently, only the extensive Raman study of Syme *et al.* (1) has been published on the vibrational spectra of thorite, but this work was performed at 295 K only. On



FIG. 1. The tetragonal unit cell of thorite in the space group D_{4h} .

the other hand, much more work (Raman and infrared studies) has been done on zircon, $ZrSiO_4$ (2-5), and hafnon, $HfSiO_4$ (6), both isomorphous with thorite.

Theoritical Considerations

Thorite has a tetragonal crystallographic cell containing four units of ThSiO₄ (Fig. 1). However, the primitive unit cell contains only two ThSiO₄ formulas and leads to 36 normal modes of vibration at k = 0. They belong to irreducible representations of the point group D_{4k} and they are listed in Table

TABLE I

Normal Modes of Thorite in the Point Group D_{4k} and Their Raman and Infrared Activities

	No	ormal r	nodes a				
D _{4h}	N	N _{ac}	N _{int}	N _{ext}		Activity	
				N _{rot}	N _{tr}	Raman	IR
A _{lg}	2	0	2	0	0	а	
A _{2g}	1	0	0	1	0		
Big	4	0	2	0	2	а	
B _{2g}	1	0	1	0	0	а	
E	5	0	2	1	2	а	
A _{lu}	1	0	1	0	0		
A _{2u}	4	1	2	0	1		a
Btu	1	0	0	1	0		
B _{2u}	2	0	2	0	0		
Eu	5	1	2	1	1		а

Note. $N_{int} + N_{ext} = N_{opt} = 2A_{1g} + 1A_{2g} + 4B_{1g} + 1B_{2g} + 5E_g + 1A_{1u} + 3A_{2u} + 1B_{1u} + 2B_{2u} + 4E_u$.

I, where N represents the total number of modes, N_{ac} the three acoustic modes, and N_{opt} the optical ones. For the optic modes the distinction between internal (N_{int}) and external (N_{ext}) modes is due to the tetrahedral substructure of the tightly bound SiO₄ molecular unit in the tetragonal structure of ThSiO₄. The internal modes involve oxygen-silicon motions within the tetrahedron and the external modes are concerned by Th⁴⁺ cation and SiO₄⁴⁻ unit motions. They can be subdivided into rotatory N_{rot} and translatory N_{tr} modes.

The Raman tensors for D_{4h} irreducible representations given by Loudon (7) are referred to the crystallographic axes X, Y, and Z as follows:

$$\mathbf{A}_{1g}\begin{bmatrix}a&\cdot&\cdot\\\cdot&a&\cdot\\\cdot&\cdot&b\end{bmatrix}\mathbf{B}_{1g}\begin{bmatrix}c&\cdot&\cdot\\\cdot&-c&\cdot\\\cdot&\cdot&\cdot\end{bmatrix}$$
$$\mathbf{B}_{2g}\begin{bmatrix}\cdot&d&\cdot\\d&\cdot&\cdot\\\cdot&\cdot&\cdot\end{bmatrix}\mathbf{E}_{g}\begin{bmatrix}\cdot&\cdot&e\\\cdot&\cdot&\cdot\\e&\cdot&\cdot\\e&\cdot&\cdot\end{bmatrix},\begin{bmatrix}\cdot&\cdot&e\\\cdot&e\\\cdot&e&\cdot\end{bmatrix}.$$

For infrared experiments, according to the selection rules, the A_{2u} and E_u modes are, respectively, allowed when the electric vector of the transmitted light is parallel and perpendicular to the fourfold crystallographic axis Z.

Experimental

ThSiO₄ exists under two crystalline forms: huttonite with a monoclinic structure and thorite with a tetragonal habit which is only concerned with this paper. The phase transition temperature occurs around 1225° C (8), the thorite being the low-temperature phase.

Crystals of thorite were grown by the flux method following Chase and Osmer (9). An equimolar mixture of ThO_2 -SiO₂ was dissolved in a convenient solvent: Li₂MoO₄-



FIG. 2. Direction of the crystallographic axes X, Y, Z with respect to the macroscopic habit of thorite.

 MoO_3 at 1150°C and slowly cooled down to around 800°C at a 1.5°C/hr rate. In these conditions, the supersaturated solution gives rise, after nucleation, to many well shaped single crystals, which do not exceed 2 mm in size.

The tetragonal structure of the obtained crystals was confirmed by X-ray analysis. The calculated cell parameters: $a = b = 7.122 \pm 0.004$ Å and $c = 6.317 \pm 0.004$ Å, are in good agreement with the literature values (10, 11). Moreover, the orientation of the crystallographic axes with respect to the tetragonal bipyramidal habit of the crystal, was checked using the Laue method (Fig. 2).

In the Raman experiments, the incident light was provided by a "Spectra Physics" cw argon laser, operating at 5145 Å. The 90° scattered light was analyzed by a "Coderg T 800" double monochromator. The incident and scattered light polarizations were determined by using Polaroid film. For an output polarization parallel to the grating lines, 100 mW of laser power was enough, but twice this power was necessary for the opposite polarization. For the low-temperature experiments, we used an "Oxford Instrument" flow cryostat equipped with a temperature regulation device between 4.2 and 300 K.

The infrared absorption experiments were performed on a "Bruker 113 V" Fourier transformed spectrometer equipped with a pyroelectric DTGC detector for the 50- to 700-cm⁻¹ region and with a MCT (Hg-Cd-Te) detector above 500 cm⁻¹. The resolution was choosen at 2 or 4 cm⁻¹ depending upon the wavenumber regions. The sample was in a crystalline powder form mixed with crude paraffin and pressed between two TPX plates for the 50to 700-cm⁻¹ region and two KBr plates above 500 cm⁻¹.

Results and Discussion

Raman

Figures 3 and 4 show the Raman spectra recorded at 300 and 4.2 K for different polarizations. The standard notation u(vw)tused for a specific geometry indicates that the incident light is directed along u and polarized along v and the scattered light is directed along t and polarized along w, where u,v,w,t are referred to the crystallographic axes X, Y, Z.

The low-temperature spectra recorded at 100 and 4.2 K (Fig. 4) are quite similar in frequencies to the room-temperature spectra (Fig. 3). The shift in energy never exceeds 2 cm⁻¹ for a same mode and the doubly degenerated E_g modes are not split at 4.2 K. In that way, we are sure that no phase transition occurs at low temperature. Therefore, the site symmetry of thorium remains D_{2d} from 300 to 4.2 K.

A lack of polarization is observed in all spectra. It should probably be due to the small size of the crystal which makes difficult the crystallographic axes orientation and provides edge scattering. For the lowtemperature experiments, the use of a win-



FIG. 3. Raman spectra of thorite recorded at room temperature for the four appropriate geometries: (1) X(ZZ)Y; (2) X(YZ)Y; (3) Z(YY)X; and (4) X(YX)Y. * Denotes a doubt in the mode assignment between B_{1g} and E_{g} symmetry.

dow cryostat increased the depolarization of the spectra. Consequently, we will consider only the room-temperature spectra for the latter discussion.

The lack of polarization is specially pronounced in spectra 3: Z(YY)X and 4: X(YX)Y in Fig. 3 which are concerned with the X and Y crystallographic axes. The same feature was noticed by Dawson *et al.* (2) who gave at that time an explanation pointing out that when these two axes are misorientated, the Raman tensors are referred to a new orthogonal set of axes, X', Y', Z', which affects only the B_{1g} and B_{2g} modes. So, their Raman tensors are transformed in such a way that both of these symmetry modes are observed together.

Because of the lack of polarization, the symmetry assignment was not obvious and we have to consider the relative intensity of each line. In particular, the three lines at 127, 597, and 919 cm^{-1} have comparable intensity for both X(YZ)Y and Z(YY)X polarizations (spectra 2 and 3 in Fig. 3); therefore they can be associated with B_{1g} and E_{g} modes. But in view of the results obtained by Syme *et al.* (1), a very weak peak at 129.5 cm^{-1} and a more intense peak at 126 cm^{-1} corresponding, respectively, to E_g and B_{1g} symmetries, one can assume that in our case the E_g mode is missing or masked by the 127 cm^{-1} which presents a halfheight width of 2.5 cm⁻¹. According to the number of each symmetry modes predicted by the group theory, these three lines should be associated with B_{1g} modes, but a doubt remains for the assignment of the 127-cm⁻¹ line.

Infrared

The analysis of the IR spectra is simpler than for Raman because only two modes



FIG. 4. Raman spectra of thorite recorded at 4.2 K for two adopted geometries: X(ZZ)Y and X(YZ)Y. * The symmetry mode could be B_{1g} .

TABLE II

PEAK FREQUENCIES ν of the Active Raman Modes of Thorite and Their Polarizations^a

v (cm ⁻¹)				
300 K	4.2 K	Polarization 300 K	Symmetry	
127	126.5	(XZ), (YZ), and (YY)	B _{1g} or	$E_{g}(e)$
194	193	(XZ), (YZ)	Eg	(e)
265		(XY)	\mathbf{B}_{2g}	(i)
294	294.5	(XZ), (YZ)	Eg	(e)
312		(YY)	\mathbf{B}_{1g}	(e)
440	439	(ZZ)	Alg	(i)
518	abs	(XZ), (YZ)	Eg	(i)
597	595.5	(XZ), (YZ) , and (YY)	Big	(i)
855	856	(XZ), (YZ)	Eg	(i)
893	894	(ZZ)	Aig	(i)
919	921	(XZ), (YZ) , and (YY)	\mathbf{B}_{1g}	(i)

^a The distinction between internal and external modes is denoted by (i) and (e).

are allowed by the selection rules: E_u and A_{2u} (see Theoretical Considerations).

The spectra recorded at room temperature for various wavenumber regions are represented in Fig. 5.

Because of the isomorphism between the structure of zircon and thorite, one can assume that the normal modes are the same in both crystals. Then the shift in frequency for each vibration in only due to the difference between the zirconium and the thorium masses. Therefore, in the absence of polarized spectra for thorite, the mode symmetries (Table III) were determined by correlation between our observed frequencies and those obtained by Dawson *et al.* (2) on the polarized infrared spectra of zircon.

General Remarks

For both infrared and Raman spectra the assignment of the internal modes was done on the basis of the free silicate measurements reported in the literature (12, 13) for a T_d symmetry group (Fig. 6). Rotatory and translatory labels for the external modes are derived from Dawson *et al.* (2).

The internal frequencies (except the B_{2g}



FIG. 5. Unpolarized infrared spectra of thorite recorded at room temperature.



FIG. 6. On the left side: Correlation diagram between the SiO₄⁴⁻ free ion vibrations in the T_d symmetry (12) and the ThSiO₄ lattice internal vibrations in the D_{4h} symmetry (where two units of SiO₄ are involved), via the D_{2d} symmetry of one SiO₄ in the crystal. On the right side: The external modes.

mode at 265 cm^{-1}) are higher in energy than the external ones, owing to the different masses engaged in each class, respectively Si and Th. However, a comparative study with the zircon frequencies points out that

TABLE III

FREQUENCIES V OF THE SEVEN ACTIVE INFRARED MODES OF THORITE AND THEIR ASSIGNMENT DEDUCED FROM THE POLARIZED INFRARED STUDY OF ZIRCON (2) ISOMORPHOUS WITH THORITE

ZrSiO₄ 300 K	ThSiO₄ 300 K	Symmetry
287	234	E _u (e)
338	332	A_{2u} (e)
389.5	356	E_{u} (e)
430.5	450	E _u (i)
608	575	$A_{2\mu}$ (i)
885	860	E _u (i)
989	911	A_{2v} (i)

the distinction between internal and external modes is not obvious. Indeed, a slight decrease of frequency from zircon to thorite for the internal motions, shows that the involved stretching vibration (Si–O) depends on the lattice environment through the influence of the oxygen-cation (Th⁴⁺ or Zr⁴⁺) force constant and unit cell volume (3). Moreover the mass effect between Th and Zr cannot explain alone the shift in frequency for the external modes. Then we have to consider that the mixing between modes of a same symmetry can modified the type of vibration given by the theoritical model.

Conclusion

The work related in this paper led us to determine all the vibration frequencies of thorite ThSiO₄ (except one Raman mode) and their symmetry in the D_{4h} factor group.

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